

QUANTUM CHEMICAL MODELING OF UV SPECTRA OF POLYURETHANE STRUCTURAL FRAGMENTS

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Results of TDDFT calculations of characteristics for excited singlet states of mono- and diisocyanates and carbamates containing from one to three phenyl groups are presented. The influence of the structural composition of the isocyanate/carbamate on the formation of its UV absorption spectrum was analyzed.

Keywords: *ab initio calculation, density functional theory, UV spectrum, isocyanates, urethanes.*

Introduction. Polyurethanes represent a modern multifaceted class of polymers with a broad range of applications in various industrial sectors owing to their ability to form macromolecular structures of defined architectures [1]. Linear segmented urethane polymers that were synthesized via the reaction of hydroxyl-containing derivatives with isocyanates included flexible and rigid segments in their structures. As a rule, the flexible segments consisted of ethers or esters and imparted to the polyurethane a high level of elasticity. The structures of the rigid short-chain segments were determined by the structures of the starting di- and polyisocyanates [2].

Experimental methods such as vibrational spectroscopy that are supplemented by results of quantum-chemical calculations are often used in research on the structural and physical properties of urethane polymers incorporating isolated functional groups in their compositions and also starting compounds for polyurethane synthesis [3–8]. Previously performed *ab initio* and DFT calculations of the structural, spectral, and energy characteristics of several urethane representatives, e.g., methylcarbamate (MC) [9], methylphenylcarbamate (MPC) [10], 4,4'-diphenylmethane-*bis*(methyl)carbamate (DPMC) [11], and a representative of the diisocyanate family and one of the polyurethanes used most frequently in industry, 4,4'-diphenylmethanediisocyanate (MDI) [12], demonstrated that the examined structural models of the polyurethane chain fragments were adequate. Nevertheless, these calculations exposed the limitations of vibrational spectroscopy methods for analyzing the structural features of the polyurethane chain because notable trends in the calculated spectra of three examined urethanes were not observed other than a tendency of the frequency of the C=O stretching vibration to decrease in the order MC → MPC → DPMC [11]. An analogous situation was typical of the vibrational spectra of phenylisocyanate (PI) and MDI [12].

Polymers containing carboxyl chromophores or amino or phenyl groups are known to absorb significantly in the UV region. However, as a rule, vibronic spectra of such polymers are unimpressive and rarely used to study their structures [13]. Nevertheless, an analysis of UV absorption bands (ABs) of polymers enables certain conclusions to be made about the number of units or functional groups in the polymer chain because the AB shifts systematically to long wavelength as the chain lengthens because the frontier molecular orbitals approach each other [13]. Time-dependent or non-stationary density functional theory (TDDFT) is currently widely used to model electronic spectra of polyatomic systems [14–20]. It was noted that a typical quantum-chemical calculation performed using the hybrid functional B3LYP in terms of the TDDFT approximation displays deviations of the excited state energies from the experimental values of the order of 0.4 eV [15]. This enables results of TDDFT calculations to be used for predictions and analysis.

Herein we present results of TDDFT quantum-chemical calculations of several characteristics of the excited electronic states of two isocyanates, PI OCN–C₆H₅ (**I**) and MDI OCN–C₆H₄–CH₂–C₆H₄–NCO (**II**), in addition to four urethanes, MC CH₃–OC(O)NH₂ (**III**), MPC CH₃–OC(O)NH–C₆H₅ (**IV**), DPMC CH₃–OC(O)NH–C₆H₄–CH₂–C₆H₄–

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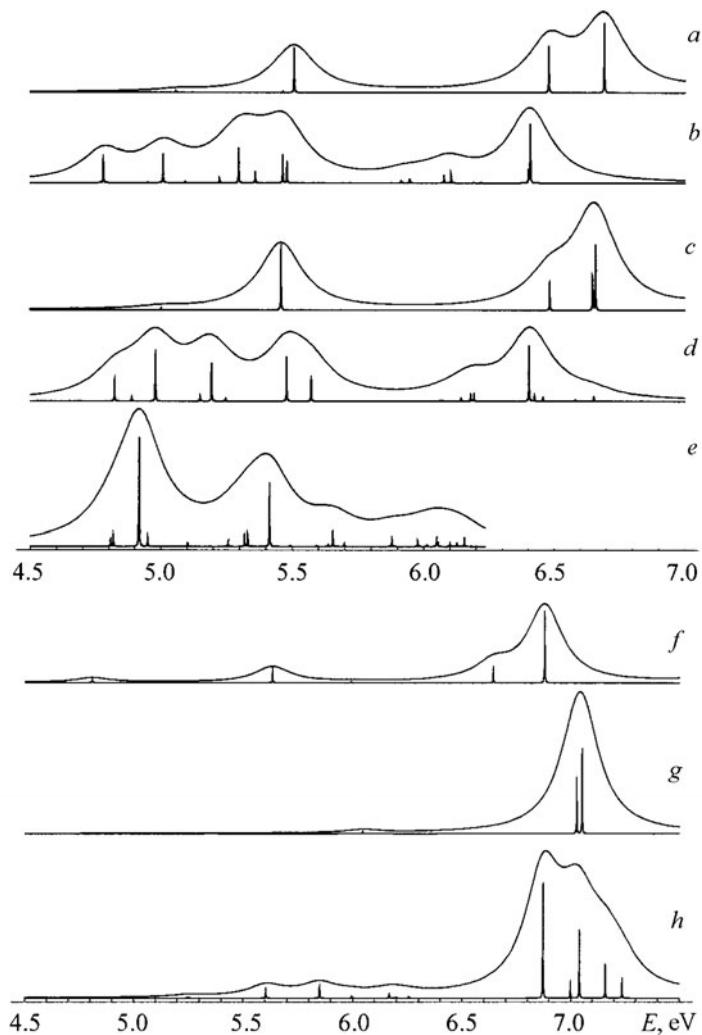


Fig. 1. Model electronic absorption spectra of **I** (a), **II** (b), **IV** (c), **V** (d), **VI** (e), **VII** (f), **VIII** (g), and **IX** (h).

$\text{HN}(\text{O})\text{CO}-\text{CH}_3$ (**V**), and a model dicarbamate **VI**, the structure of which was based on the DPMC chain with the incorporation into the central part of the molecule of a single phenyl and a single methyl: $\text{CH}_3-\text{OC(O)NH-C}_6\text{H}_4-\text{CH}_2-\text{C}_6\text{H}_4-\text{CH}_2-\text{C}_6\text{H}_4-\text{HN}(\text{O})\text{CO}-\text{CH}_3$. These urethanes include the principal functional groups found in a typical polyurethane (urethane, phenyl, methyl, and methylene) and can act as models of polymer structural fragments.

Calculation Procedure. Spectral characteristics of the examined compounds were calculated using the applied quantum-chemical program GAMESS [21, 22]. The results were visualized using the program MacMolPlt [23]. All calculations were performed in the standard basis set cc-pVDZ [24] using TDDFT methods and the hybrid exchange-correlated functional B3LYP [25–27]. The equilibrium configurations of **I–IV** that were found earlier [9, 10, 12] were used to model their absorption spectra. Full geometry optimization in the B3LYP/cc-pVDZ approximation based on the equilibrium structure of DMPC obtained earlier in the HF/6-311G approximation [11] was carried out for **V** and **VI**.

Results and Discussion. Energy characteristics of 20 excited singlet electronic states were calculated for each of **I–V**; 30 states, for **VI**. Table 1 presents the results from calculations for the 20 lowest states. Figure 1 shows model spectra of **I**, **II**, **IV**, and **V** at excitation energies up to 7 eV and of **VI** at excitation energies up to 6.3 eV. Systems of energy states of aniline $\text{C}_6\text{H}_5-\text{NH}_2$ (**VII**), toluene $\text{C}_6\text{H}_5-\text{CH}_3$ (**VIII**), and diphenylmethane $\text{C}_6\text{H}_5-\text{CH}_2-\text{C}_6\text{H}_5$ (**IX**) were also calculated in order to assign more completely the ABs to chromophores. Figure 1 shows model spectra of these compounds at excitation energies up to 7.5 eV.

TABLE 1. Energies of Lower Excited Singlets of Electron States, Wavelengths, and Transition Oscillator Strengths from the Ground State

I			II			III			IV			V			VI		
<i>E</i> , eV	λ , nm	<i>S</i>	<i>E</i> , eV	λ , nm	<i>S</i>	<i>E</i> , eV	λ , nm	<i>S</i>	<i>E</i> , eV	λ , nm	<i>S</i>	<i>E</i> , eV	λ , nm	<i>S</i>	<i>E</i> , eV	λ , nm	<i>S</i>
5.06	245	0.012	4.78	259	0.158	6.85	181	0	5.00	248	0.014	4.82	257	0.125	4.81	258	0.040
5.46	227	0.007	4.95	251	0.003	7.55	164	0.007	5.46	227	0.366	4.89	254	0.028	4.82	257	0.069
5.51	225	0.250	5.01	248	0.166	7.93	156	0.034	6.20	200	0	4.98	249	0.286	4.92	252	0.602
6.43	193	0	5.09	243	0.012	8.35	149	0.040	6.43	193	0	5.15	241	0.041	4.95	250	0.074
6.48	191	0.263	5.22	237	0.031	8.50	146	0.072	6.48	191	0.167	5.19	239	0.215	5.03	247	0.002
6.58	188	0.002	5.30	234	0.198	8.66	143	0.149	6.64	187	0.187	5.25	236	0.019	5.10	243	0.022
6.69	185	0.389	5.35	232	0.002	8.80	141	0.031	6.67	186	0.361	5.48	226	0.251	5.26	236	0.033
6.96	178	0	5.36	231	0.068	8.88	140	0.011	7.01	177	0.001	5.57	223	0.143	5.32	233	0.068
7.20	172	0	5.46	227	0.162	8.98	138	0.024	7.15	173	0.054	6.06	205	0.006	5.33	232	0.069
7.46	166	0.352	5.48	226	0.125	9.11	136	0.026	7.38	168	0.001	6.07	204	0.006	5.35	231	0.007
7.49	165	0.007	5.91	210	0.018	9.43	132	0.052	7.42	167	0.015	6.14	202	0.022	5.41	230	0.354
7.78	160	0.107	5.95	209	0.025	9.68	128	0.015	7.46	166	0.002	6.18	201	0.047	5.50	226	0.006
7.80	159	0.001	6.08	204	0.046	9.86	126	0.013	7.49	166	0.029	6.19	200	0.047	5.60	222	0.004
7.86	158	0	6.10	203	0.056	9.87	126	0.006	7.61	163	0.019	6.22	199	0.001	5.64	220	0.013
7.88	157	0	6.19	200	0.002	10.03	124	0.021	7.70	161	0	6.31	196	0.002	5.66	219	0.087
7.98	155	0	6.22	199	0.002	10.19	122	0.008	7.88	158	0.003	6.40	194	0.310	5.70	218	0.017
8.05	154	0.004	6.40	194	0.069	10.20	122	0.003	7.92	157	0.001	6.42	193	0.044	5.80	214	0
8.18	152	0.001	6.41	194	0.330	10.22	121	0.016	7.98	156	0.006	6.46	192	0.024	5.88	211	0.058
8.32	149	0.002	6.43	193	0	10.24	121	0.016	8.00	155	0.035	6.58	188	0.010	5.98	207	0.035
8.40	148	0.002	6.44	193	0	10.31	120	0.059	8.01	155	0.019	6.65	186	0.032	5.98	207	0
$\Delta E_{\text{H-L}}$, eV																	
5.74			5.33			8.54			5.82			5.44			5.36		

Let us first analyze model electronic spectra of compounds containing only one phenyl (**I** and **IV**), noting that their structures are exceedingly complicated (Fig. 1a and c). ABs were observed at 275, 230, and 222 nm in the experimental vibronic spectrum of PI in the range 300–200 nm [28]. The first of these (much weaker) demonstrated clearly resolved vibrational structure analogous to the vibronic transition in the spectrum of benzene at ~260 nm [29]. Thus, the weak AB at 245 nm (5.06 eV) in the spectrum of **I**; 248 (5.01) (**IV**), 258 (4.81) (**VII**), and 232 (5.34) (**VIII**) corresponded to the forbidden transition $^1A_{1g} \rightarrow ^1B_{2u}$ (256 nm) in the spectrum of benzene [30]. Bands at 193 and 185 nm (6.43 and 6.69 eV) in the spectrum of **I**; at 193 and 187 (6.43 and 6.64) in that of **IV**; at 191 and 180 (6.49 and 6.89) in that of **VII**; and at 205 and 175 (6.05 and 7.09) in the spectrum of **VIII** corresponded with the two other electronic transitions of benzene that fell in the UV region below 170 nm {also forbidden $^1A_{1g} \rightarrow ^1B_{1u}$ (203 nm) and allowed $^1A_{1g} \rightarrow ^1E_{1u}$ (180) [30]}.

Two transitions in the range 200–230 nm (6.20–5.39 eV), one of which (longer wavelength in spectra of **IV** and **VII**) was rather strong; the other, close to zero, corresponded to absorption of the aniline amine (**VII**), the MPC NH (**IV**), and the PI –N=C= N-containing fragment (**I**). Toluene (**VIII**) did not absorb in this range in the calculated spectrum. It is noteworthy that the calculation in general reproduced adequately the system of PI lower singlet states. The energy of the first excited state was elevated by 0.55 eV; of the second, by 0.07; of the third, lowered by 0.08.

Absorption of the C=O chromophore [30] in spectra of **I** and **IV** corresponded to a strong AB at 191 nm (6.48 eV). Absorption of the methyl was seen as strong ABs at 186 and 176 nm (6.67 and 7.05 eV) in spectra of **IV**

and **VIII**. It should be mentioned that such assignments were somewhat arbitrary because up to 30 pairs of molecular orbitals contributed to these transitions as a result of strong intramolecular interaction.

Electronic spectra of compounds incorporating into their structures two benzene rings and a methylene group (**IX**) in addition to two isocyanates (**II**) or urethanes (**V**) exhibited complicated structures (Fig. 1b and d). The most characteristic difference of the spectra of these compounds and those examined above was a significant bathochromic shift of the ABs. Furthermore, the symmetry of the molecular fragments was reduced (this refers mainly to the benzene rings) together with an overall increase in the number of chromophores, which increased the number of ABs. In turn, this was accompanied by a strengthening of previously forbidden transitions. An example was the two strong transitions in the range 260–240 nm (4.77–5.17 eV) in spectra of **II** and **V** that were analogs of the aforementioned $^1A_{1g} \rightarrow ^1B_{2u}$ transition in the spectrum of benzene. Weak ABs due to the presence of a methylene group in the structures of **II**, **V**, and **IX** also appeared in this range. The features mentioned above (bathochromic shift and strengthening of ABs) were not observed in the spectrum of diphenylmethane (**IX**). This enabled them to be explained mainly as a manifestation of intramolecular interaction (inductive effect). The short-wavelength spectral range 200–175 nm (6.20–7.09 eV) also underwent substantial changes. Only one of two strong ABs remained at 194 nm (6.29 eV). It was due primarily to absorption of the carbonyl chromophore (Fig. 1a–d).

Adding additional moieties (benzene and methylene groups) to the DMPC structure did not cause a significant bathochromic shift of the ABs in the spectrum of **VI** relative to the position of the ABs in the spectrum of **V** (Table 1). Nevertheless, such a structural change was accompanied by a redistribution of intensity in the long-wavelength region because the transition at 252 nm (4.92 eV) due to absorption of the central benzene ring that was missing in the structures of **II** and **IV** dominated.

The spectrum of the simplest urethane **III** differed in principle from those of the other urethanes and isocyanates because it was determined by absorption of amine and urethane chromophores. The ABs in the spectrum of **III** experienced a significant hypsochromic shift relative to that in spectra of **IV** and **V**. Thus, the carbonyl absorption appeared at 146 nm (8.50 eV); amine, 143 (8.66). The energy differences of the highest occupied and lowest unoccupied molecular orbitals ΔE_{H-L} (Table 1) in the series **I** → **II** and **III** → **IV** → **V** → **VI** demonstrated a smooth decrease caused by the increasing number of functional groups in their structures.

Conclusion. Changes of electronic absorption spectra that were due to different numbers of functional groups in the structures of isocyanate or urethane chains were analyzed using quantum-chemical TDDFT modeling of electronic structures of isocyanates and carbamates containing one or two isocyanate or urethane groups in addition to one to three benzene rings. It was shown that the position of the ABs in the electronic spectrum was determined mainly by the number of structural elements (phenyl and urethane/isocyanate groups) incorporated into the carbamate/isocyanate. The influence of the terminal functional groups (urethane or isocyanate) on the formation of the electronic spectrum structure seemed less significant. The established features could be used for analytical purposes.

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