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# RAMAN SCATTERING IN DIAMOND IRRADIATED WITH HIGH-ENERGY XENON IONS

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**Abstract.** Accumulation of radiation damage and associated mechanical stresses in diamonds irradiated with 167 MeV Xe ions to fluences  $1.0 \times 10^{10} \div 8.15 \times 10^{14} \text{ cm}^{-2}$  have been studied using confocal Raman spectroscopy. The spectra were measured in a backscattering geometry across the irradiated layer by scanning the edge of single crystalline synthetic samples with nitrogen concentration  $3 \div 5 \text{ ppm}$ . All spectra were recorded at room temperature. Parameters of the  $1332 \text{ cm}^{-1}$  first-order Raman line – the FWHM and peak position, studied as function of Xe ion fluence, were used to characterize the structural disorder and mechanical stress level.

**Key words:** diamond, ion implantation, radiation-induced defects, amorphization, stress, Raman scattering

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## **1. Introduction**

Raman excitation in diamond is usually accompanied by photoluminescence (PL) and both Raman and PL spectra parameters are very sensitive to structural damage [1,2]. For instance, critical dose of amorphization of ion irradiated diamond can be estimated by due to appearance of  $1530\text{ cm}^{-1}$  broad band [1,3–5]. Raman scattering on vibrations of intrinsic defects at  $1490$  and  $1630\text{ cm}^{-1}$  provides information on vacancy and interstitial defects, respectively [1,4]. Highly sensitive PL methods can distinguish individual impurity and defect centers [2,6], that can be used to get impurity/defect distribution in the irradiated layer [7]. As known, any radiation damage of diamond lattice leads to its swelling [8–14]. Structural disorder induced by low energy ions (tens to hundreds keV) is localized near the surface. The boundary between the damaged layer and the undamaged part of the crystal is mechanically strained. For the high energy implantation (tens to hundreds MeV) the boundary of damaged layer is located beyond tens of micrometres from the surface due to radiation damage produced in elastic collisions in the end of the ion range. Crystal lattice, in this case, is inhomogeneously expanded and strained along ion range. Earlier studies of the depth distribution of impurities/defects in diamond crystals irradiated by high energy ions [7] were done using precision removal of thin layers of superhard material and utilizing PL the technique with fixed excitation and registration parameters [15]. Modern micro Raman spectrometers, combined with a 3D scanning confocal microscope and equipped with several laser excitation sources, allow performing such work at the new level. The aim of this work is depth profiling of radiation damage and mechanical stresses in diamond, irradiated with 167 MeV Xe ions.

## **2. Experimental**

The samples used in this research were a few millimetre size plates cut along the (100) plane from single crystal type Ib synthetic diamonds produced by “AdamasInvest”, Belarus by High Pressure High Temperature (HPHT) method [16]. After mechanical polishing of both sides the thickness of the specimen was about  $800\text{ }\mu\text{m}$ . Average concentration of nitrogen in the diamonds plates was calculated from visible and UV absorption spectra [16]. Visible and

UV absorption spectra were measured at room temperature in spectral range 200–900 nm using spectrometer Cary 300 Bio (Agilent Technologies). The size of the measured areas was about 0.6 mm in diameter.

Irradiation with 167MeV Xe ions was carried out at room temperature to fluences in the range  $1.0 \times 10^{10} - 8.15 \times 10^{14} \text{ cm}^{-2}$  at the IC-100 FLNR, JINR cyclotron [17]. After getting of definite fluence, the specimen was taken out from the irradiation module and one fourths of it was screened by a thick copper tape. Thus, four areas with different ion fluences were created on one diamond plate.

Raman spectra were registered using a micro Raman spectrometer Nanofinder High End (LOTIS TII Japan-Belarus) combined with a 3D scanning confocal microscope. Spectral resolution of the spectrometer was  $2.3 \text{ cm}^{-1}$  and  $0.77 \text{ cm}^{-1}$  for 600 lines/mm and 1800 lines/mm gratings correspondingly. Maximum spectral resolution of  $0.25 \text{ cm}^{-1}$  was reached with Echelle grating (75 lines/mm). Excitation was provided by blue (473 nm) and green (532 nm) lasers with powers of 4 and 10 mW respectively and a spot size of 1  $\mu\text{m}$ . Raman spectra were gathered using a backscattering setup in the range  $50\text{-}5700 \text{ cm}^{-1}$ .

A schematic view of the experimental setup is shown in Fig.1a. A photo of a polished cross-section of the diamond plate for the fluence  $8.15 \times 10^{14} \text{ cm}^{-2}$  with a focal point of a green laser is shown in Fig.1b. Raman spectra were registered on polished diamond cross-sections along ion trajectory (along X axis) by scanning sample with a deepening of the excitation region (Z coordinate) by 1.5 or 5.0  $\mu\text{m}$ . Scanning was started beyond the sample at a distance of about 5  $\mu\text{m}$  from the edge. The mid-point of the rising curve of the main Raman line intensity was chosen as the origin of the X-coordinate. Scanning step was 0.5 or 1  $\mu\text{m}$ . Exposure time varied from 10 to 60 seconds and depended on the Raman signal intensity. Regions with low nitrogen concentration (3-5 ppm) were selected for examination. Such cross section geometry guaranteed the measurement of Raman spectra in the irradiated layer, corresponding to the level of damage at a selected depth and provided a higher spatial resolution.

Energy loss and depth profiles of carbon vacancies and implanted Xe ions were calculated using Monte Carlo method [18] with SRIM 2000 software [19]. Calculation results are shown in Fig.2. For 167 MeV Xe, the projected ion range in diamond is  $R_p = 10.67 \mu\text{m}$  with  $0.44 \mu\text{m}$  straggling. Carbon vacancy concentration calculated for single ion varied from  $2.24 \times 10^6 \text{ cm}^{-3}$  at the surface region to  $6.66 \times 10^7 \text{ cm}^{-3}$  at the area of maximum defects creation (at the depth of  $10.45 \mu\text{m}$ ). According to calculations no vacancies were created beyond the projected ion range at the depth more than  $11 \mu\text{m}$ .

### 3. Results and discussion

The Raman spectrum of the virgin diamond specimen has a narrow band at  $1332 \text{ cm}^{-1}$  originated from single-phonon light scattering and well defined two-phonon scattering band in the range from  $2200$  to  $2660 \text{ cm}^{-1}$ . The full width at half maximum (FWHM) of the main line was about  $1.6 \text{ cm}^{-1}$ . This value of the Raman line width is typical for structurally perfect type IIa diamonds [20–23].

The evolution of the Raman/PL spectra in the surface layer with increasing ion fluence is shown in Fig. 3. Noticeable changes occurred after the fluence more than  $1 \times 10^{11} \text{ cm}^{-2}$ . The intensity of the main Raman line decreased, narrow lines at  $3530 \text{ cm}^{-1}$ ,  $3120 \text{ cm}^{-1}$ ,  $1635 \text{ cm}^{-1}$ ,  $1425 \text{ cm}^{-1}$  and wide bands in the vicinity of  $1230 \text{ cm}^{-1}$  and  $400 \text{ cm}^{-1}$  appeared. The intensity of lines peaked at  $3530 \text{ cm}^{-1}$ ,  $3120 \text{ cm}^{-1}$ ,  $1425 \text{ cm}^{-1}$  initially increased (approximately up to a fluence  $6 \times 10^{13} \text{ cm}^{-2}$ ), and then decreased with increasing fluence. The intensity of broad bands  $1220 \text{ cm}^{-1}$  and  $400 \text{ cm}^{-1}$  and a narrow line  $1635 \text{ cm}^{-1}$  grew monotonically. At fluence  $6 \times 10^{13} \text{ cm}^{-2}$  in the vicinity of the  $1635 \text{ cm}^{-1}$  line, Raman/PL spectrum changed and a characteristic series of narrow lines in the region of  $1400\text{--}1900 \text{ cm}^{-1}$  appeared.

Narrow lines at  $3530 \text{ cm}^{-1}$  ( $743 \text{ nm}$ ),  $3120 \text{ cm}^{-1}$  ( $638 \text{ nm}$ ), и  $1425 \text{ cm}^{-1}$  ( $575 \text{ nm}$ ) are ascribed to radiative recombination [2]. The  $743 \text{ nm}$  line (PL center GR1) is associated with vacancies in a neutral charge state. The lines  $638 \text{ nm}$  (PL center  $\text{NV}^-$ ) and  $575 \text{ nm}$  (PL center

$NV^0$ ) are associated with the nitrogen – vacancy complex in the negatively charged and neutral charge states, respectively.

Wide bands in the vicinity of  $400\text{ cm}^{-1}$  and  $1230\text{ cm}^{-1}$  and series of narrow lines in the area of  $1400\text{--}1900\text{ cm}^{-1}$  refers to Raman scattering in diamond. Wide bands at  $400\text{ cm}^{-1}$  and  $1230\text{ cm}^{-1}$  were observed in Raman spectra of diamonds irradiated by ions [3,4,24–26] and neutrons [26,27]. It is assumed that they are associated with amorphous  $sp^3$  bonded carbon clusters along the ion tracks i.e. clusters of carbon displaying short-range  $sp^3$  ordering without any long range order. These clusters are localized under high pressure inside the substantially intact diamond matrix and, therefore, are not transformed into  $sp^2$ -bonded carbon structure. Series of narrow lines in the spectral region  $1400\text{--}1900\text{ cm}^{-1}$  was associated with Raman scattering on intrinsic point defects in a disordered by irradiation diamond lattice [28]. The most intense lines at  $1500$  and  $1635\text{ cm}^{-1}$  are related [1,3,4], respectively, to vacancies and intrinsic interstitials. The peak positions of these lines shift to lower wavenumbers and FWHM are increased almost linearly with increasing damage [4].

The evolution of Raman spectra with depth in the irradiated layer is shown in Fig. 4a and 4b for two Xe ion fluences. As can be clearly seen that for  $6.1 \times 10^{12}\text{ cm}^{-2}$ , the intensity of the radiation induced PL centers GR1,  $NV^-$ ,  $NV^0$  decreased with depth, demonstrating correlation with nuclear stopping profile. The PL bands due to complex defects  $NV^-$ ,  $NV^0$  completely disappeared at a depth of about  $8\text{ }\mu\text{m}$ , while the simple vacancy center GR1 was registered up to the end of ion range. Intrinsic interstitial related Raman line at  $1635\text{ cm}^{-1}$  had the similar behavior - both GR1 and  $1635\text{ cm}^{-1}$  lines were observed up to a depth of  $10.5\text{ }\mu\text{m}$ . The intensity of the broad Raman bands at  $400\text{ cm}^{-1}$  and  $1230\text{ cm}^{-1}$  increased from the surface into the depth of the irradiated layer synchronously with the growth of radiation damage concentration. Raman/PL spectra beyond the ion range reverted to their initial state before irradiation and composed of the only main Raman line  $1332\text{ cm}^{-1}$ .

For the fluence  $8.15 \times 10^{14} \text{ cm}^{-2}$  PL was substantially suppressed even in the near surface layer (see Fig.3), therefore, the evolution of the spectra shown in Fig. 4b was due to Raman scattering. The increase of radiation damage in the diamond lattice with depth caused a decrease in intensity, broadening and low frequency shift of all Raman lines. At the region of the maximum disorder of diamond lattice (at a depth of about 9-10  $\mu\text{m}$ ), the main Raman line  $1332 \text{ cm}^{-1}$  disappeared. The spectrum consisted only of broad bands at  $400 \text{ cm}^{-1}$  and  $1220 \text{ cm}^{-1}$ , associated with amorphous carbon clusters, and a wide band appeared at  $1530 \text{ cm}^{-1}$  (Fig. 4b). A low-intensity line at  $1635 \text{ cm}^{-1}$  from interstitial carbon was registered on the high-frequency wing of the  $1530 \text{ cm}^{-1}$  band. The first appearance of the  $1530 \text{ cm}^{-1}$  band in the Raman spectra corresponded to fluence of  $3.55 \times 10^{14} \text{ cm}^{-2}$  when it was registered in the damage peak region. All Raman bands ascribed to radiation defects disappeared at a depth more than 10  $\mu\text{m}$  and the only scattering from the intact diamond crystal lattice is detected ( $1332 \text{ cm}^{-1}$  line).

The  $1530 \text{ cm}^{-1}$  band was observed in diamonds irradiated with different ions [1,3–5,9,14,29–31]. Its appearance was associated with the critical density of radiation damage in diamond, above which the graphitization of the damaged layer occurred during annealing [32]. Spectral position and shape of the  $1530 \text{ cm}^{-1}$  band coincided with those of highly disordered by ion irradiation graphite  $\text{sp}^2$  layers [1,5,33–35]. The formation of amorphous, not graphite, carbon as a result of irradiation of diamond with heavy ions was confirmed by direct measurements of electron diffraction [36].

The absence of the  $1332 \text{ cm}^{-1}$  line in the Raman spectra at a depth of 9–10  $\mu\text{m}$  was an evidence of the practically complete destruction of the crystalline  $\text{sp}^3$  phase of diamond. The appearance of the  $1530 \text{ cm}^{-1}$  band instead of the  $1332 \text{ cm}^{-1}$  line, which is associated with the  $\text{sp}^2$  phase, was a sign of the transition of diamond to the state of amorphous carbon. The amorphized layer was not yet continuous, but consisted of islands, since at a depth of amorphization, a line at  $1635 \text{ cm}^{-1}$  from intrinsic internodes was registered. As shown in [14], Raman scattering from defects in diamond is absent when complete amorphization occurs. Due

to high sensitivity of the Raman technique to  $sp^2$ -bonded carbon structures, the absence of a wide band  $1530\text{ cm}^{-1}$  was convincing evidence that implantation of Xe ions with the energy 167 MeV did not lead to formation of amorphous carbon for irradiation fluences less than  $3.55 \times 10^{14}\text{ cm}^{-2}$ .

According to the intensity distribution of  $400\text{ cm}^{-1}$  and  $1230\text{ cm}^{-1}$  bands, amorphous  $sp^3$ -bonded carbon clusters formed along the Xe ion trajectory even at relatively low fluences. At the same time, due to the lack of detailed studies of ion-irradiated diamond with transmission electron microscopy (TEM), there is no direct evidence confirming the existence of  $sp^3$ -coupled clusters in diamond irradiated with swift heavy ions.

Figures 5a and 5b show variation in spectral position and FWHM of main Raman line  $1332\text{ cm}^{-1}$  with Xe ion fluence. The gap between curves, plotted for the maximum fluences (the results of two independent measurements on two samples are shown), corresponded to the depth of the irradiated layer, where the  $1332\text{ cm}^{-1}$  line was absent in the Raman spectra, but a wide band of  $1530\text{ cm}^{-1}$  was registered. For all fluences, the low-frequency shift and FWHM of the  $1332\text{ cm}^{-1}$  line grew monotonically with a depth up to  $9.5\text{ }\mu\text{m}$  synchronously with calculated (see Fig. 2) nuclear energy losses profile. More significant electronic energy losses, as in [37], did not manifest themselves in defect formation. The same dependence was noted by all researchers of ion-irradiated diamonds (see, for example, [4,14,38–40]). It should be noted, that the depth of maximum defect formation determined experimentally from the shift and FWHM of the  $1332\text{ cm}^{-1}$  line was approximately  $1.0\text{ }\mu\text{m}$  closer to the surface than predicted by SRIM -  $10.45\text{ }\mu\text{m}$ . Similar differences in the defect formation depths between calculated and experimental values in ion irradiated diamond samples were observed also in [4].

The inhomogeneous defects distribution caused elastic stresses in both the irradiated layer and the interface region with intact diamond as follows from the shift and broadening of the  $1332\text{ cm}^{-1}$  line. An increase of ion fluence leads to monotonic increment of the spectral shift and FWHM values of the  $1332\text{ cm}^{-1}$  line, but, in general, did not affect their depth profiles in

the irradiated layer (see Fig. 5). Change rate of the spectral shift and FWHM of main Raman line with increasing fluence was greater in the end of range region than in subsurface layer. The monotonic increment of shift and FWHM with an increase of ion fluence was interrupted after  $3.55 \times 10^{14} \text{ cm}^{-2}$  when the change in the parameters of the  $1332 \text{ cm}^{-1}$  line seemed to be saturated. At the boundary with the incipient amorphous layer, the maximum shift of the  $1332 \text{ cm}^{-1}$  line was about  $20 \text{ cm}^{-1}$ , and its maximum broadening was about  $60 \text{ cm}^{-1}$ . In the region of projected ion range this “saturation” might be associated with the beginning of formation of carbon amorphous phase islands. Accordingly to SRIM, number vacancies for fluence  $3.55 \times 10^{14} \text{ cm}^{-2}$  is about  $2.36 \cdot 10^{22} \text{ cm}^{-3}$  in the damage peak region, that is comparable with critical threshold of radiation damage in diamond [14,32,35,36,41]. The number of accumulated vacancies in the near-surface region for this fluence is almost 30 less and the  $1332 \text{ cm}^{-1}$  line was clearly detected, but its position and FWHM remained practically unchanged with a further accumulation of radiation damage. The maximum shift of the  $1332 \text{ cm}^{-1}$  line in subsurface region was less than  $5 \text{ cm}^{-1}$ , and its maximum broadening was about  $20 \text{ cm}^{-1}$ .

The absence of changes in the near-surface parameters distribution of the  $1332 \text{ cm}^{-1}$  line with the onset of amorphization may be due to partial relaxation of stresses in the entire irradiated layer after formation of an amorphous layer and a loss of bond integrity between the irradiated and undamaged layers in the diamond crystal lattice. Similar effects were previously described for silicon implanted with 2 MeV Xe ions [42] and 1.5 MeV Si ions [43], in which the formation of amorphous layer was accompanied by stress relaxation. One can suggest that increase in interatomic distances in the irradiated layer was not restrained by anything along the normal to the irradiated surface. Lattice swelling in direction perpendicular to the irradiated surface was observed as an additional increment in the thickness of the irradiated layer [11–14].

As known, mechanical stresses induced by the tangential expansion of the irradiated layer may result in bending even of the whole diamond plate [12,14,42]. The low-frequency

shift of the main Raman line  $1332\text{ cm}^{-1}$ , corresponding to the expansion of the diamond crystal lattice, is linearly connected with elastic stresses [24,44–46]. According to [46] we estimated the magnitude of the elastic stresses  $\sigma$  in the  $\langle 001 \rangle$  direction using relation:

$$\sigma = 0.49 \cdot [\text{GPa} / \text{cm}^{-1}] \cdot \Delta\nu [\text{cm}^{-1}],$$

where  $\Delta\nu$  is the shift of the spectral position of the  $1332\text{ cm}^{-1}$  line relative to that of an intact diamond sample. For instance, taking into account the above expression, the magnitude of elastic stresses for the fluence  $3.55 \times 10^{14}\text{ cm}^{-2}$  is about 2.5 GPa in subsurface layer and about 10 GPa in damage peak region.

Figure 6 shows the behaviour of the  $1332\text{ cm}^{-1}$  line parameters beyond the irradiated layer up to the opposite side of the diamond plate. Raman spectra were measured with a resolution of  $0.25\text{ cm}^{-1}$  and a deepening of the excitation region on  $5\text{ }\mu\text{m}$ . Spectral position ( $1331.9\text{ cm}^{-1}$ ) and FWHM ( $1.59\text{ cm}^{-1}$ ) for unirradiated sample are shown by a dashed line. As can be seen from Fig. 6a, the spectral position of the main Raman line beyond the irradiated layer restored nonmonotonically to its value before irradiation. For relatively small fluences (not more than  $1 \times 10^{13}\text{ cm}^{-2}$ ), the lattice expansion corresponding to the low-frequency shift of the  $1332\text{ cm}^{-1}$  line gradually decreased with depth, reaching its original values at a depth of about  $60\text{--}70\text{ }\mu\text{m}$ . Stress distribution beyond the ion range for high fluences was nearly identical: the expansion of the lattice was changed by its compression at a depth of about  $3\text{ Rp}$ , followed by an increase in compression stresses up to  $6\text{--}7\text{ Rp}$ . Further, with increasing depth, compression stresses monotonously weakened, but did not reach the value before irradiation even on the opposite side of the plate (at a depth of about  $800\text{ }\mu\text{m}$ ).

As follows from Figure 6, the unirradiated part of the diamond plate “restrained” the expansion of the crystal lattice in the irradiated layer. With this in mind, it is expected that the stress distribution in diamonds irradiated with swift heavy ions will depend on the thickness of the plates.

## Conclusions

Accumulation of radiation damage and associated mechanical stresses in diamond samples irradiated by 167 MeV Xe ions were studied with Raman spectroscopy. The main Raman line  $1332\text{ cm}^{-1}$  became less intense, broadened and shifted to lower frequencies region due to radiation damage of the diamond lattice. Stress distribution profiles in irradiated layer and beyond it were derived from spectral position of the  $1332\text{ cm}^{-1}$  line. The expansion of the lattice in the irradiated by 167 MeV Xe ions,  $10\text{ }\mu\text{m}$  layer leads to appearing tangential stresses that cause deformation (bending) of the entire diamond wafer. The undamaged part of the diamond lattice is under alternating tensile and compressive stresses that "compensate" for the expansion of the lattice in the irradiated layer. Partial relaxation of mechanical stresses was observed in the entire irradiated layer, indicating its lower hardness and associated with the onset of amorphization in the region of maximum defect concentration.

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Figure captions:

**Fig. 1.** Schematic view of the Raman analysis measurements (a) and a photo of a polished cross-section of the diamond plate for the fluence  $8.15 \times 10^{14} \text{ cm}^{-2}$  (b).

**Fig. 2.** Energy loss profiles, primary carbon vacancies and implanted impurity distributions for Xe ions with the energy 167 MeV. Results of calculations are given for the fluence  $1 \text{ cm}^{-2}$ .

**Fig. 3.** Raman/PL spectra evolution in the near surface layer ( $z = 1.5 \text{ }\mu\text{m}$ ) with increasing irradiation fluence of Xe ions with the energy 167 MeV. Values of fluence [ $\text{ions}/\text{cm}^2$ ] are indicated in labels to curves. Spectra are shown without displacement along the ordinate axis.

**Fig. 4.** Raman/PL spectra evolution with the irradiated layer depth for ion fluences (a)  $6.1 \times 10^{12} \text{ cm}^{-2}$  and (b)  $8.15 \times 10^{14} \text{ cm}^{-2}$ . Spectra are shown without displacement along the ordinate axis.

**Fig. 5.** Depth profile of (a) spectral position and (b) FWHM of the  $1332 \text{ cm}^{-1}$  Raman line in diamond samples irradiated with 167 MeV Xe ions for different ion fluences.

**Fig. 6.** FWHM (a) and spectral position (b) depth profiles of  $1332 \text{ cm}^{-1}$  line beyond the ion range in diamond plate irradiated with 167 MeV Xe ions for different ion fluences. Dashed line corresponds to virgin diamond sample.











