

# DETERMINATION OF THE PHONON AMPLITUDES EMPLOYED IN BOSON EXPANSION THEORIES

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Low-lying states linked up by strong enough  $E2$ -transitions can be interpreted as superposition of many collective phonon configurations  $(D^+)^{nd}$ . In Boson Expansion Theories these phonons are mapped onto quadrupole boson operators, e.g. in [1]  $D^+ \rightarrow d^+ \sqrt{1 - \hat{n}_d / \Omega} \sim d^+ s$ ,  $\hat{n}_d$  is the number operator of  $d$ -boson,  $\Omega$  is their maximum number,  $s$  is a scalar auxiliary boson of the Interacting Boson Model. Wave functions in such approach are found as solutions of the boson Hamiltonian ( $H_B$ ). Theoretical calculations of  $H_B$  parameters begin with a choice of the  $D$ -phonon operator for which we use the form of the Quasiparticle Random Phase Approximation (QRPA):  $D_\mu^+ \sim \sum (\psi_{12} a_1^+ a_2^+ + \varphi_{12} a_{2^-} a_{1^-})$  where amplitudes  $\psi, \varphi$  determine its two-quasiparticle composition. However, in contrast to QRPA where  $\psi, \varphi$  are defined for a one-phonon state, we search them taking into account many phonon (boson) structure of the collective states. With this object we minimize over  $\psi$  and  $\varphi$  a functional  $\Phi$  comprising the expectation value of  $H_B$  (parameters of which are functions of  $\psi, \varphi$  and effective quasiparticle interactions) and some additional conditions, details can be found in [2]. Thus, equations for  $\psi, \varphi$  involve the boson expectation values (BEV) such as  $\langle n_d \rangle$ ,  $\langle d^+ \cdot d^+ s s + \text{H.c.} \rangle$  and others, i.e. the equations allow for the many boson structure of states. One of the addition conditions in  $\Phi$  fixes the value of  $\xi = \sum \varphi^2 / \sum \psi^2$  to be  $\ll 1$ , that gives the possibility, first, to employ the usual QRPA calculations and, secondly, to obtain the selfconsistent description of  $\psi, \varphi$  and BEV, i.e. the  $H_B$  parameters calculated with final values of  $\psi, \varphi$  give such BEV which being substituted into equations for  $\psi, \varphi$  lead to the same values of parameters. Such selfconsistency is impossible in the Tamm-Dankoff method, i.e., when  $\varphi = 0$ . Calculations for Xe isotopes have shown that  $\xi$  cannot be larger than 0.065 and a reasonable agreement between calculations and all experimental data can be attained with  $0.012 < \xi \leq 0.050$ . A part of these calculations for <sup>122</sup>Xe is given in the table ( $E$  in MeV, МэВ,  $B(E2)$  in  $e^2 \text{fm}^4$ ).

	$E(2_1^+)$	$E(2_2^+)$	$E(4_1^+)$	$B(E2: 2_1^+ \rightarrow 0_1^+)$	$B(E2: 4_1^+ \rightarrow 2_1^+)$
Exp.	0.331	0.843	0.828	$2890_{-165}^{+125}$	$4150_{-180}^{+190}$
$\xi = 0.0145$	0.330	0.838	0.882	2920	4360
$\xi = 0.02$	0.339	0.843	0.894	2860	4270
$\xi = 0.03$	0.329	0.858	0.866	2800	4210

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