

DETERMINATION OF THE PHONON AMPLITUDES EMPLOYED IN BOSON EXPANSION THEORIES

Efimov A.D.¹, Jolos R.V.², Mikhajlov V.M.³

¹*Ioffe Physical Technical Institute, St.-Petersburg, Russia;* ²*Joint Institute for Nuclear Research, Dubna, Russia;* ³*Physical Institute of St.-Petersburg State University, Russia*

E-mail: efimov98@mail.ru

Low-lying states linked up by strong enough $E2$ -transitions can be interpreted as superposition of many collective phonon configurations $(D^+)^{nd}$. In Boson Expansion Theories these phonons are mapped onto quadrupole boson operators, e.g. in [1] $D^+ \rightarrow d^+ \sqrt{1 - \hat{n}_d} / \Omega \sim d^+ s$, \hat{n}_d is the number operator of d -boson, Ω is their maximum number, s is a scalar auxiliary boson of the Interacting Boson Model. Wave functions in such approach are found as solutions of the boson Hamiltonian (H_B). Theoretical calculations of H_B parameters begin with a choice of the D -phonon operator for which we use the form of the Quasiparticle Random Phase Approximation (QRPA): $D_\mu^+ \sim \sum (\psi_{12} a_1^+ a_2^+ + \varphi_{12} a_{2-} a_{1-})$ where amplitudes ψ, φ determine its two-quasiparticle composition. However, in contrast to QRPA where ψ, φ are defined for a one-phonon state, we search them taking into account many phonon (boson) structure of the collective states. With this object we minimize over ψ and φ a functional Φ comprising the expectation value of H_B (parameters of which are functions of ψ, φ and effective quasiparticle interactions) and some additional conditions, details can be found in [2]. Thus, equations for ψ, φ involve the boson expectation values (BEV) such as $\langle n_d \rangle$, $\langle d^+ \cdot d^+ s s + \text{H.c.} \rangle$ and others, i.e. the equations allow for the many boson structure of states. One of the addition conditions in Φ fixes the value of $\xi = \sum \varphi^2 / \sum \psi^2$ to be $\ll 1$, that gives the possibility, first, to employ the usual QRPA calculations and, secondly, to obtain the selfconsistent description of ψ, φ and BEV, i.e. the H_B parameters calculated with final values of ψ, φ give such BEV which being substituted into equations for ψ, φ lead to the same values of parameters. Such selfconsistency is impossible in the Tamm-Dankoff method, i.e., when $\varphi \equiv 0$. Calculations for Xe isotopes have shown that ξ cannot be larger than 0.065 and a reasonable agreement between calculations and all experimental data can be attained with $0.012 < \xi \leq 0.050$. A part of these calculations for ^{122}Xe is given in the table (E in MeV, $M\text{eB}$, $B(E2)$ in e^2fm^4).

| | $E(2_1^+)$ | $E(2_2^+)$ | $E(4_1^+)$ | $B(E2: 2_1^+ \rightarrow 0_1^+)$ | $B(E2: 4_1^+ \rightarrow 2_1^+)$ |
|----------------|------------|------------|------------|----------------------------------|----------------------------------|
| Exp. | 0.331 | 0.843 | 0.828 | 2890_{-165}^{+125} | 4150_{-180}^{+190} |
| $\xi = 0.0145$ | 0.330 | 0.838 | 0.882 | 2920 | 4360 |
| $\xi = 0.02$ | 0.339 | 0.843 | 0.894 | 2860 | 4270 |
| $\xi = 0.03$ | 0.329 | 0.858 | 0.866 | 2800 | 4210 |

1. Р.В.Джолос, Ф.Дэнау, Д.Янссен // ТМФ. 1974. Т.20. С.112; F.Donau, D.Janssen, R.V.Jolos // Nucl. Phys. A. 1974. V.224. P.93.

2. A.D.Efimov, V.M.Mikhajlov // Bull. RAS. Ser. Phys. 2011. V.75. P.890.