



# Photoluminescence and enhanced chemical reactivity of amorphous SiO<sub>2</sub> films irradiated with high fluencies of 133-MeV Xe ions



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## ABSTRACT

SiO<sub>2</sub>/Si structures have been irradiated with 133 MeV Xe<sup>+17</sup> ions at fluencies of  $(10^{10}-5 \times 10^{14}) \text{ cm}^{-2}$ . The structure transformation and light-emitting properties of irradiated SiO<sub>2</sub> films were studied using RS, SEM, TEM and PL techniques as well as chemical etching in 4% solution of hydrofluoric acid (HF). An intensive photoluminescence in visible range was registered from the samples irradiated at a fluence of  $10^{14} \text{ cm}^{-2}$  and higher. Simultaneously, it was found a drastic increase of SiO<sub>2</sub> etch velocity in HF solution for the irradiated samples. Annealing (1100 °C, 2 h) of irradiated samples resulted in PL quenching and etch velocity recovery practically to the value of non-irradiated SiO<sub>2</sub>. It was concluded that radiative oxygen deficient centers are responsible for the PL appearance. It was also shown that the etch velocity ratio of the irradiated and virgin SiO<sub>2</sub> in 4%-HF ( $V_{\text{irr}}/V_{\text{virgin}}$ ) can be used in order to estimate the radiation damage in SiO<sub>2</sub> matrix irradiated with high fluencies of swift heavy ions.

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## 1. Introduction

Swift heavy ion (SHI) projectiles can produce point defects, defect clusters and ion latent tracks – extended damage regions along ion trajectories - in an irradiated target. Ion tracks were first observed in nuclear fission reactors [1] and since then they have been widely studied due to their growing list of technological applications. It is possible to modify thin films of insulators by means of SHI irradiation. Thin SiO<sub>2</sub> films integrated into silicon wafers are of special interest for nanotechnology. The passage of swift ions through SiO<sub>2</sub> and some other insulators is usually described by the thermal spike model [2–5]. In this model, it is assumed that latent tracks are the result of interaction of the projectile ions with the target electrons, and an initial track size is determined by a cylindrical molten zone (future latent track) created along the ion trajectory. Owing to a different chemical reactivity than virgin matrix, the latent track regions in SiO<sub>2</sub> can be transformed into nano-

microchannels by means of etching in hydrofluoric acid (HF) solutions [6–9]. The other promising field is a possibility for controlling the shape and size of nanoclusters embedded in dielectric matrix via SHI irradiation. The authors of Refs [10–13] have reported a shape transformation of Ag, Au, Co and InAs nanoparticles embedded into amorphous SiO<sub>2</sub> films from spherical one to conical or elongated along the SHI beam direction. A new wave of interest in studying the SHI–SiO<sub>2</sub> interaction is due to a possible usage of SHIs for the creation of Si-based light-emitting nanostructures via SiO<sub>2</sub> disproportionation in the ion tracks [14,15].

Moreover, when solar cells and MOS structures with insulating silica layers are in outer space equipment, the radiation defects generation in SiO<sub>2</sub> is possible. The point defects and defect clusters play a very important role in the silica's electronic properties and could modify the electronic device's characteristics. Thus, radiation induced defects in silica deserve a special attention in the modern electronic technology, too.

Latent track regions in SiO<sub>2</sub> are referred to as radiation damaged areas. In etching solution, the damaged areas of irradiated sample are etched much faster than the undamaged ones [16]. It provides an opportunity to transform the tracks into pores by means of

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controlled chemical etching. Though, this possibility could be realized only for low SHIs fluencies. Multiple tracks overlapping takes place at high fluencies, and the *a*-SiO<sub>2</sub> layer, on the whole, becomes damaged. Hence, the etch velocity of irradiated SiO<sub>2</sub> matrix ( $V_{irr}$ ) in HF solution should be higher in comparison with the etch velocity of virgin one ( $V_{virgin}$ ). In our study, we have used this parameter ( $V_{irr}$ ) to estimate the SHI-induced damage in SiO<sub>2</sub> matrix.

The aim of this paper is to investigate chemical reactivity and light-emitting properties of SiO<sub>2</sub> layers irradiated with high SHIs fluencies. We have also tried to elucidate a correlation between SHI-induced photoluminescence (PL) and structural transformations of SiO<sub>2</sub> as well as to estimate the influence of high-temperature treatment on the etch velocity and the photoluminescence of irradiated SiO<sub>2</sub> layers.

## 2. Methods

The  $1 \times 1 \text{ cm}^2$  samples were cut from the thermally oxidized Si (100) wafer. The thickness of SiO<sub>2</sub> film was  $\approx 1 \mu\text{m}$  according to the ellipsometry data. These samples were afterwards irradiated with 133 MeV Xe<sup>17+</sup> ions to fluencies ranging from  $1 \times 10^{10}$  to  $5 \times 10^{14} \text{ cm}^{-2}$  at room temperature and at the normal incidence. Irradiation was conducted at the DC-60 cyclotron (Astana, Kazakhstan). After irradiation, a part of the samples was additionally annealed at 1100 °C for 2 h in Ar atmosphere.

The electronic energy loss  $S_e$  and the nuclear energy loss  $S_n$  for SiO<sub>2</sub>/Si structure were simulated using the SRIM'2013 code [17]. We have calculated the diameter and lifetime of molten regions along the Xe ion trajectories in SiO<sub>2</sub> in the frame of inelastic thermal spike model [5]. It enables us to estimate a latent track size as well as to calculate an extent of latent track overlapping for our experimental conditions. The SiO<sub>2</sub> density was taken to be  $2.2 \text{ g/cm}^3$  in computer simulations [18].

The surface morphology and microstructure were analyzed by scanning electron microscopy (SEM) and transmission electron microscopy in “cross-section” technique (XTEM). The scanning microscopy study was performed using a Hitachi S-4800 microscope at an electron-beam orientation perpendicular to the surface. The microstructure of irradiated SiO<sub>2</sub>/Si samples was investigated by a 200-keV Hitachi H-800 microscope. The standard procedure of sample's mechanical polishing and thinning by ion milling was used in TEM cross-sections preparation. Raman scattering (RS) and photoluminescence (PL) were used to study optical properties of the samples. Raman spectra were measured using the micro-Raman setup Ntegra Spectra with a laser beam at  $\lambda = 473 \text{ nm}$ . PL signal was excited by a He–Cd laser source at  $\lambda = 325 \text{ nm}$ . All optical measurements were performed at room temperature.

The virgin and irradiated samples as well as the irradiated samples after annealing were treated in a 4% aqueous solution of hydrofluoric acid (HF) at room temperature for 6 min. A half of the surface of each sample was covered with wax prior etching. In order to avoid artefacts all the samples were etched in the same process (simultaneously). The thickness of SiO<sub>2</sub> layer for protected and unprotected regions of the etched samples was evaluated from laser ellipsometry ( $\lambda = 632.8 \text{ nm}$ ). Thus, we could estimate the thickness of SiO<sub>2</sub> layer removed during the HF-solution treatment. The etch velocity could be determined as a ratio of the removed layer thickness to the etching time.

## 3. Results and discussion

According to the SRIM code,  $S_e$  in the SiO<sub>2</sub> layer is nearly constant and amounts to  $\sim 13.5 \text{ keV/nm}$ . This value is about 2 orders of magnitude higher than  $S_n$  in SiO<sub>2</sub>. The calculated radius of molten region (future latent track) formed along the 133 MeV Xe ion

trajectory in SiO<sub>2</sub> and its lifetime were 5.9 nm and 21.5 ps, respectively [5]. The molten region lifetime allows us speaking about the presence of the liquid phase in SiO<sub>2</sub> matrix and its subsequent solidification.

The question is if the calculated radius of molten region corresponds to the radius of resulting latent track. It is well known that direct observation of ion tracks in *a*-SiO<sub>2</sub> is a very difficult procedure due to a lack of sufficient contrast between the amorphous matrix and the track (see, for example [19]). That is why, a number of indirect methods such as IR absorption [16], small angle X-ray scattering (SAXS) [20] and track etching technique [21,22] were used successfully to quantify the sensitivity of amorphous SiO<sub>2</sub> under swift heavy ion irradiation. In the recent review [19], the results gathered by each of these techniques were compared with each other to gain a better understanding of track formation phenomena. The latent track radii deduced from IR spectroscopy, SAXS analyses as well as extrapolated ones from the etching experiments were plotted as a function of  $S_e$  for different specific beam energies ranged from 0.14 to 11 MeV/u. The  $S_e$  of Xe ions with 1 MeV/u in SiO<sub>2</sub> (our experimental conditions) is equal to  $\sim 13.5 \text{ keV/nm}$ . According to the “radius versus  $S_e$ ” dependence from Ref. [19] (obtained from SAXS data for Au ions at 0.14–0.94 MeV/u), the latent track radius amounts to  $\sim 5.5 \text{ nm}$  at the 13.5 keV/nm. It is in a reasonable agreement with the calculated radius of molten region (5.9 nm). Hence, we can use the calculated value for preliminary estimation of the extent of latent track overlapping.

One can estimate the extent of latent track overlapping using the equation

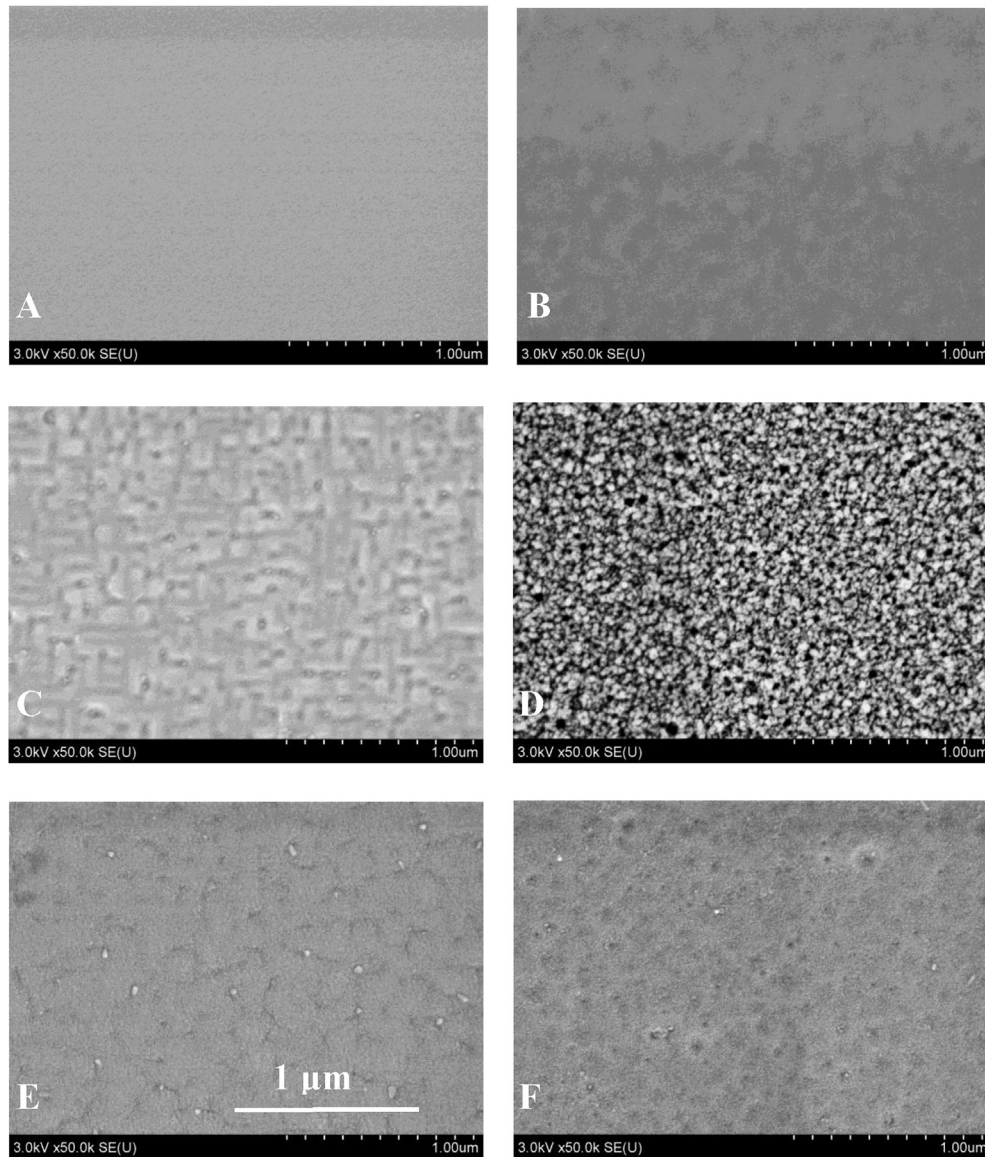
$$K = \pi R_{tr}^2 \Phi,$$

where  $K$  denotes the extent of overlapping,  $R$  is the calculated radius of molten region and  $\Phi$  is the ion fluence. The calculated  $K$  values are summarized in Table 1. One can see that track overlapping is negligible at the fluence of  $1 \times 10^{10} \text{ cm}^{-2}$ . The irradiated area of silica is completely covered with tracks at the fluence of  $1 \times 10^{12} \text{ cm}^{-2}$ . The further fluence increase results in multiple overlapping.

The surface morphology features of the samples under investigation were examined by SEM. Fig. 1 presents the SEM surface micrographs of the virgin SiO<sub>2</sub>/Si and the samples irradiated to fluencies of  $1 \times 10^{12} \text{ cm}^{-2}$  and  $5 \times 10^{14} \text{ cm}^{-2}$  after etching in HF-solution. The surface of virgin SiO<sub>2</sub> layer is without any peculiarities (not shown), and etching does not result in any surface morphology changes (Fig. 1A). As it was mentioned above, a part of the surface of each sample was protected during the etch procedure and was not subjected to an influence of etch solution. Hence, the surface morphology of these regions is the same as for the as-irradiated SiO<sub>2</sub>/Si samples before etching. Further, we will denote protected surfaces of the samples as the as-irradiated ones. Surprisingly, the pattern formation has been observed at the as-irradiated surfaces (Fig. 1C, E). Worm-like features have been evolved on the surface of the samples irradiated to the fluence of  $1 \times 10^{12} \text{ cm}^{-2}$  (Fig. 1C). These features have got less pronounced (more coarse), and additional features of bright contrast (hillocks)

**Table 1**  
The extent of track overlapping ( $K$ ) for fluencies used in this experiment.

Fluence, $\text{cm}^{-2}$	$K$
$1 \times 10^{10}$	$\sim 0$
$1 \times 10^{12}$	1
$1 \times 10^{14}$	109
$5 \times 10^{14}$	547



**Fig. 1.** The surface of  $\text{SiO}_2/\text{Si}$  samples after etching in 4% HF for 6 min: the virgin (without irradiation) sample (A), the sample irradiated with 133 MeV  $\text{Xe}^{17+}$  ions to the fluence of  $5 \times 10^{14} \text{ cm}^{-2}$  and annealed at  $1100 \text{ }^\circ\text{C}$  for 2 h (B), the sample irradiated with 133 MeV  $\text{Xe}^{17+}$  ions to the fluence of  $1 \times 10^{12} \text{ cm}^{-2}$  (C, D) and the sample irradiated with 133 MeV  $\text{Xe}^{17+}$  ions to the fluence of  $5 \times 10^{14} \text{ cm}^{-2}$  (E, F). The sample surface was protected during the etch procedure (C, E). The scale is the same for all micrographs.

have appeared with fluence increasing to  $5 \times 10^{14} \text{ cm}^{-2}$  (Fig. 1E). A formation of self-organized surface patterns on a silica glass substrate due to high energy ion irradiation has been recently reported in Ref. [23]. As in our experiment, worm-like patterns on the surface of silica samples normally irradiated with 1.8 MeV  $\text{Au}^{2+}$  ions to the fluence higher than  $5 \times 10^{16} \text{ cm}^{-2}$  were observed, too. The authors of Ref. [23] have explained the observed surface modification via compressive stresses generating in the near surface region during the irradiation. In our opinion, mechanisms of the pattern formation on the surface of ion-irradiated materials are complicated and not yet understood. A clearance of the origin of surface modification of irradiated samples is a subject of future investigations and beyond the scope of this paper.

After irradiation the samples were etched in 4% HF solution as described in Sect. 2. The etched surface of the sample irradiated to the fluence of  $1 \times 10^{12} \text{ cm}^{-2}$  is characterized by overlapping irregular shape pores (Fig. 1D). Further fluence increasing results in a substantial surface transformation. Only a few shallow pits are

observed on the etched surface of the samples irradiated to fluences of  $1 \times 10^{14}$  (not shown) and  $5 \times 10^{14} \text{ cm}^{-2}$  (Fig. 1F). The track overlapping values amount to 109 and 547 at  $1 \times 10^{14}$  and  $5 \times 10^{14} \text{ cm}^{-2}$ , respectively (see Table 1). What is the reason for such inhomogeneous etching despite the complete damage overlapping at  $10^{12} \text{ cm}^{-2}$  and multiple one at higher fluences? This effect can be explained by a complicated structure of individual tracks in a- $\text{SiO}_2$  demonstrated by Kluth et al. [20] recently using SAXS measurements. They revealed a track fine structure comprising a cylindrical core-shell configuration with a lower density core and a higher density shell as compared to un-irradiated material. More recently, Kluth et al. [24] reported a steady state of nanoscale density fluctuations in a- $\text{SiO}_2$  irradiated with high fluencies of SHIs when the surface is completely covered by ion tracks. SAXS measurements and molecular dynamics simulations were consistent that these density fluctuations resulted from track overlapping where newly formed tracks annihilated (partially) pre-existing tracks. Apparently, well-developed

morphology (at  $10^{12} \text{ cm}^{-2}$ ) and shallow pits (at higher fluencies), we have observed on the etched surfaces of irradiated samples are the signature of density fluctuations discussed in Ref. [24].

Fig. 1B depicts the view of the etched surface of the sample irradiated to the fluence of  $5 \times 10^{14} \text{ cm}^{-2}$  and afterwards annealed at  $1100 \text{ }^\circ\text{C}$  for 2 h. One can see substantial smoothing of the etched surface of the annealed sample. A similar effect (smoothing of the etched surface) has been observed for all irradiated and annealed samples (not shown). A disappearance of the as-irradiated surfaces patterns is observed after annealing, too. In consequence, the surface images of the etched and as-irradiated areas after annealing become the similar ones. This is not surprising as the thermal treatment leads to radiation damage recovery as it will be shown below from the etch velocity and PL data.

Fig. 2 presents the etch velocity in 4% HF solution for irradiated and annealed  $\text{SiO}_2$  films as a function of ion fluence.  $V_{\text{virgin}}$  of non-irradiated sample amounts to  $17.0 \text{ nm/min}$  and is marked as a point at the vertical axis. The etch velocity of irradiated samples increases with the fluence and reaches the highest value ( $70.2 \text{ nm/min}$ ) at the fluence of  $5 \times 10^{14} \text{ cm}^{-2}$ . One can see that the heat treatment at  $1100 \text{ }^\circ\text{C}$  for 2 h results in the etch velocity recovery practically to its virgin value, obviously, due to damage annealing in the irradiated a- $\text{SiO}_2$ .

For explanation of observed results, at first, let us discuss a mechanism of latent track transformation into a pore during chemical etching. A rearrangement of energetically favorable six-component rings consisting of  $\text{SiO}_4$ -tetrahedra into three- and four-component rings takes place along the latent SHI track in a- $\text{SiO}_2$  [16,25]. The increase of the number of small rings (with a decreased angle of the Si–O–Si bonds) leads to a decrease of the amorphous silicon dioxide volume [25,26]. This is a reason for the emergence of microscopic defects with stressed Si–O bonds in the form of paramagnetic  $E'$  centers [27]. Oxygen deficient defects emerge at the damaged areas by means of the knock-on process of oxygen atoms [28]. The presence of  $E'$  centers and oxygen deficient centers with strained bonds around them explains the higher etch velocity along the SHI track [27]. The difference of etch velocities between track region and undamaged  $\text{SiO}_2$  matrix, in turn, results in the pore formation on the spot of latent track. Such scenario is realized for low SHI fluencies. Though, if the track density is high enough, at high fluencies, individual track regions are merged, and the SHI-induced microscopic defects with stressed Si–O bonds

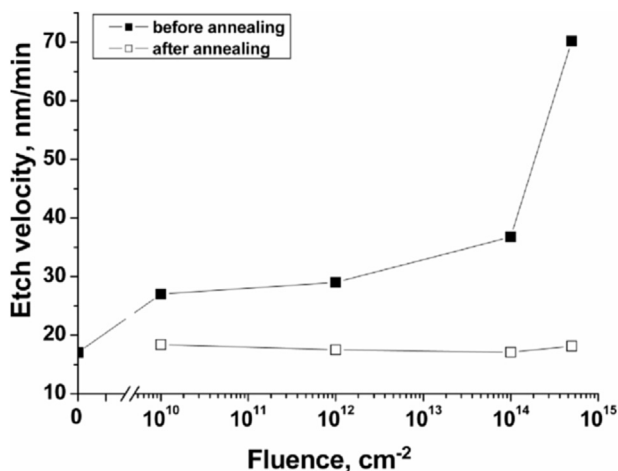


Fig. 2. The etch velocity of  $\text{SiO}_2$  films irradiated at different fluencies: (—■—) as-irradiated samples, (—□—) irradiated samples after annealing at  $1100 \text{ }^\circ\text{C}$  for 2 h (the lines are there to guide the eyes). The etch velocity of virgin sample is marked as a point at the vertical axis.

around them are distributed all over the  $\text{SiO}_2$  matrix. We can suggest an enhanced etching rate of such material in comparison with virgin one. In our experiment, track overlapping begins at the fluence of  $1 \times 10^{12} \text{ cm}^{-2}$  (see Table 1). Though, there is no substantial difference of the  $\text{SiO}_2$  etch velocities between the samples irradiated to the fluence of  $1 \times 10^{10}$  or  $1 \times 10^{12} \text{ cm}^{-2}$ . For comparison, the  $V_{\text{irr}}/V_{\text{virgin}}$  ratio amounts to 1.6 and 1.7 for fluencies of  $1 \times 10^{10}$  and to  $1 \times 10^{12} \text{ cm}^{-2}$ , respectively. At  $1 \times 10^{14} \text{ cm}^{-2}$ , multiple track overlapping takes place, and the number of microscopic defects increases substantially. The etch velocity of irradiated  $\text{SiO}_2$  increases, too. At the fluence of  $5 \times 10^{14} \text{ cm}^{-2}$  a further accommodation of defects responsible for enhanced etching rate results in the etch velocity drastic increasing. The  $V_{\text{irr}}/V_{\text{virgin}}$  ratio amounts to 2.2 and 4.1 for fluencies of  $1 \times 10^{14} \text{ cm}^{-2}$  and  $5 \times 10^{14} \text{ cm}^{-2}$ , respectively. Let us compare this ratio with the similar parameter  $V_T/V_B$  of the individual track etching process in a- $\text{SiO}_2$ . In regarding to individual track etching,  $V_B$  is the etch velocity of the undamaged material, whereas  $V_T$  is the etch velocity of damaged material along the ion track. The  $V_T/V_B$  ratio is usually in the range of  $2 < V_T/V_B < 4$  [9]. One can see that the  $V_{\text{irr}}/V_{\text{virgin}}$  ratio under multiple track overlapping conditions (obtained in our experiment) is in the same range. In other words, if the  $V_T/V_B$  ratio is a measure of track region damage at low SHI fluencies, then the  $V_{\text{irr}}/V_{\text{virgin}}$  ratio can be used in order to estimate the radiation damage in the  $\text{SiO}_2$  matrix irradiated with high SHI fluencies.

Fig. 3 shows the PL and reflectance spectra for the virgin and irradiated samples. The PL spectrum of the virgin sample includes one low-intensive band centered at 450 nm. No additional features are found after the irradiation with the fluence of  $1 \times 10^{10} \text{ cm}^{-2}$ . A new band centered at 550 nm appears in the spectrum of the sample irradiated with the fluence of  $1 \times 10^{12} \text{ cm}^{-2}$ . Fluence increasing to  $1 \times 10^{14} \text{ cm}^{-2}$  results in significant growth of PL intensity (by 10 times) and in appearance of the band set with maxima at 390 nm (ultra-violet), 450 nm (blue), 550 nm (green, the most intensive one) and 670 nm (red). The further fluence increase to  $5 \times 10^{14} \text{ cm}^{-2}$  leads to the intensity redistribution between blue-violet and red bands and to a negligible red shift of green band in comparison with the sample irradiated at the fluence of  $1 \times 10^{14} \text{ cm}^{-2}$ . We believe that these four PL bands cannot be

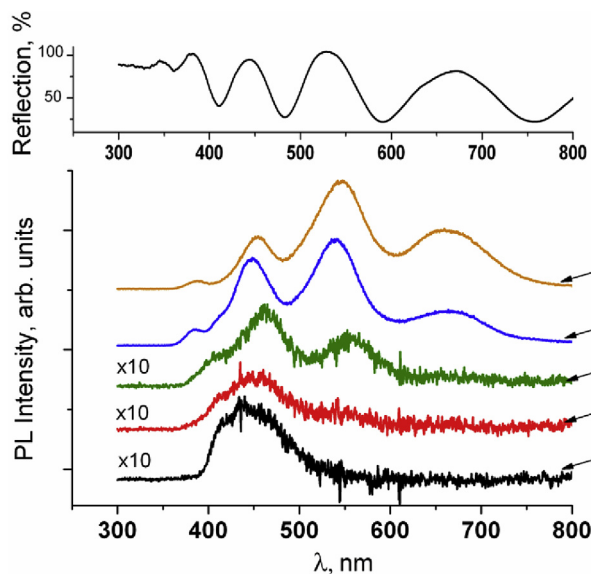


Fig. 3. PL spectra of  $\text{SiO}_2$  films: virgin (1), after irradiation with 133 MeV Xe ions to the fluence of  $1 \times 10^{10} \text{ cm}^{-2}$  (2),  $1 \times 10^{12} \text{ cm}^{-2}$  (3),  $1 \times 10^{14} \text{ cm}^{-2}$  (4) and  $5 \times 10^{14} \text{ cm}^{-2}$  (5). The upper spectrum depicts the reflectance of the sample irradiated to the fluence of  $1 \times 10^{14} \text{ cm}^{-2}$ .

attributed to different light-emitting sources. More likely, their appearance is caused by an interference effect. In this case, the thickness ( $d$ ) of the film can be estimated from the PL spectra using the formula

$$d = \frac{\lambda_1 \lambda_2}{2n(\lambda_2 - \lambda_1)},$$

where  $n=1.46$  is the average refractive index of the film in the spectral range under investigation, and  $\lambda_1$  and  $\lambda_2$  are the wavelengths of consecutive peaks. The calculated SiO<sub>2</sub> film thickness is equal to 900 nm. This is a rough estimation because SiO<sub>2</sub> film refractive index is not constant through the spectral range of the measurements. It can be expected that SHI-irradiation should change the refractive index value, too. Though, the calculated SiO<sub>2</sub> thickness is correlated with the SiO<sub>2</sub> film thickness measured by ellipsometry (1  $\mu$ m). To prove our assumption of the interference effect, the reflectance spectrum was also registered (Fig. 3). The similar oscillation trend and coincidence of maxima in PL and reflectance spectra confirm the light interference after multiple reflections within the film. Taking the interference effect into account, one can suppose that the true PL of the samples irradiated with high fluencies ( $1 \times 10^{14}$  and  $5 \times 10^{14}$  cm<sup>-2</sup>) could be depicted as one white broad band centered at 550 nm. Annealing does not result in a noticeable PL spectrum change of the virgin sample and of the sample irradiated at  $1 \times 10^{10}$  cm<sup>-2</sup>. On the contrary, PL quenching after annealing has revealed for the samples irradiated to higher fluencies ( $1 \times 10^{12}$ – $5 \times 10^{14}$  cm<sup>-2</sup>). A weak band at 440 nm is only recorded for these samples (not shown).

A passage of SHI through the silica can be followed by the formation of light-emitting Si nanoclusters in the ion tracks due to SiO<sub>2</sub> disproportionation. Structural and optical properties of SiO<sub>2</sub>/Si structures irradiated with 167 MeV Xe ions and with 700 MeV Bi ions in a fluence range of  $10^{12}$ – $10^{14}$  cm<sup>-2</sup> have been studied in Refs [14,15], and an appearance of yellow-orange photoluminescence after irradiation has been reported. PL intensity increased with the ion fluence up to  $10^{14}$  cm<sup>-2</sup> in the case of Xe and saturated at  $5 \times 10^{12}$  cm<sup>-2</sup> in the case of Bi. The electronic energy loss amounts to  $\approx 14.5$  keV/nm for 167 MeV-Xe ions. This value is higher in comparison with  $S_e \approx 13.5$  keV/nm for 132 MeV-Xe ions (in our experiment). Though, a difference in 1.0 keV/nm is unsubstantial. It allows us to compare our results with data from Refs. [14,15] obtained for a-SiO<sub>2</sub> irradiated with Xe ions. In Refs [14,15] SHI-induced PL was ascribed to the formation of Si-rich nanostructures inside the ion tracks. This suggestion has been proved by scattering at  $480$  cm<sup>-1</sup> registered in RS spectra of the irradiated samples. Its feature is typical of atomic vibrations in the amorphous Si.

For our samples, on the contrary to these results, RS method did not reveal any signs of the formation of amorphous or nanocrystalline Si phase in the SiO<sub>2</sub> samples neither after irradiation nor after annealing (not shown). The intensive peak at  $520$  cm<sup>-1</sup> belonging to the crystalline Si wafer dominated in all RS spectra. We did not register the amorphous Si band at  $480$  cm<sup>-1</sup> in irradiated samples. Commonly, the furnace annealing with a high thermal budget is applied [29–31] for the crystallization of amorphous nanoclusters and for a nanocluster's size increase. One could expect in our case that long-time annealing at  $1100$  °C should stimulate the crystallization of Si nanoclusters created in the ion tracks due to SiO<sub>2</sub> disproportionation (if any!). A presence of silicon nanocrystals should result in an appearance of low-energy tail of the Si-band at  $520$  cm<sup>-1</sup> in RS spectra [32,33]. Though, we did not observe any Si-band shape distortion in RS spectra of irradiated and annealed SiO<sub>2</sub>/Si structures.

The next step to reveal the formation of Si inclusions in SiO<sub>2</sub> was

XTEM investigation of the irradiated samples. XTEM revealed no any structural transformation or precipitation in as-irradiated SiO<sub>2</sub> films. XTEM images of irradiated SiO<sub>2</sub> films after annealing were without any structural peculiarities, too (not shown). Hence, TEM results correlate with RS data and prove a conclusion of the Si nanoprecipitates absence in the irradiated SiO<sub>2</sub> films.

Thus, we can conclude that PL of the SiO<sub>2</sub> layers irradiated with swift Xe ions is not related with the Si nanoclusters formation. The other possible explanation of visible emission is a defect generation in the SiO<sub>2</sub> matrix by SHI irradiation. At first, let us discuss the luminescence of as-deposited SiO<sub>2</sub> and SiO<sub>2</sub> irradiated to the lowest fluence ( $1 \times 10^{10}$  cm<sup>-2</sup>). The PL spectrum of virgin SiO<sub>2</sub>/Si sample includes only one weak band at 450 nm. A similar PL band was observed for the pure SiO<sub>2</sub> in Ref. [34] and attributed to F<sub>2</sub> colour centre. An existence of such colour centres is due to the formation of oxygen deficient defects. The irradiation with Xe ions to the fluence of  $1 \times 10^{10}$  cm<sup>-2</sup> does not result in any noticeable PL spectrum transformation. As it is mentioned above the latent tracks are the regions of high concentration of oxygen deficient defects. Though, track overlapping for the fluence of  $1 \times 10^{10}$  cm<sup>-2</sup> is negligible. Correspondingly, a part of irradiated SiO<sub>2</sub> matrix occupied with the latent tracks is negligible, too, and a damage level of SiO<sub>2</sub> matrix, as a whole, remains low enough. Annealing does not alter the PL spectra of the virgin sample and the sample irradiated to the lowest fluence. Taking this into account, we attribute the band at 450 nm to native defects in silicon dioxide.

A weak green band at 550 nm in the PL spectrum of the sample irradiated to the fluence of  $1 \times 10^{12}$  cm<sup>-2</sup> can be ascribed to F<sub>4</sub> oxygen deficient center [34]. Obviously, the oxygen deficient defect concentration increases with fluence, and the formation of more complicated radiative centers such as F<sub>4</sub> takes place in the irradiated layer. As it was shown above, the true PL of the samples irradiated to the higher fluencies ( $\geq 10^{14}$  cm<sup>-2</sup>) could be depicted as one white broad band centered at 550 nm. The similar white band was observed in the PL spectra of a-SiO<sub>2</sub> irradiated with 5 MeV Kr ions to fluencies of ( $10^{15}$ – $10^{16}$ ) cm<sup>-2</sup> [34]. This PL band was attributed to the formation of radiative oxygen deficient centers. Under condition of multiple tracks overlapping, we can suggest the formation of a variety of oxygen deficient defect complexes in irradiated SiO<sub>2</sub>. Apparently, these features are responsible for white PL of the irradiated a-SiO<sub>2</sub> observed in our experiment.

Fig. 4 depicts the PL peak intensity as a fluence function. One can

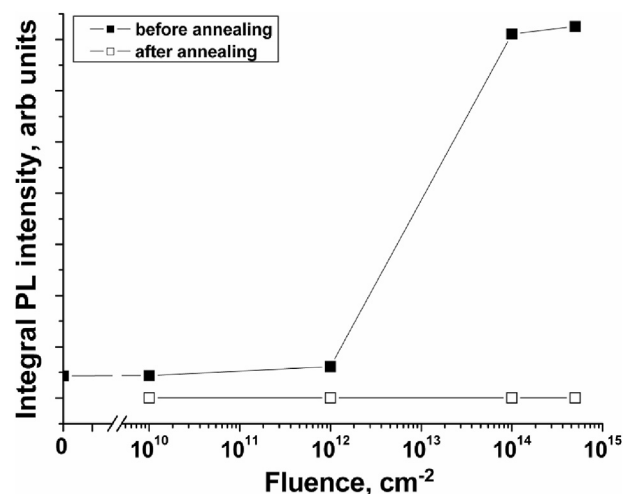


Fig. 4. The integral PL intensity of SiO<sub>2</sub> films irradiated at different fluencies: (—■—) as-irradiated samples, (—□—) irradiated samples after annealing at  $1100$  °C for 2 h (the lines are there to guide the eyes). The PL intensity of virgin sample is marked as a point at the vertical axis.

see that in the range of low fluencies ( $10^{10}$  -  $10^{12}$ )  $\text{cm}^{-2}$  the PL signal intensity is weak and practically the same as one from the non-irradiated sample. The PL signal increases by 10 times with the fluence increase from  $10^{12}$  to  $10^{14}$   $\text{cm}^{-2}$  and saturates at  $5 \times 10^{14}$   $\text{cm}^{-2}$ . The authors of Ref. [15] reported “the PL peak intensity-fluence” dependence for a-SiO<sub>2</sub> films irradiated with 167 MeV Xe ions in a fluence range of  $10^{12}$ – $10^{14}$   $\text{cm}^{-2}$ . One can see from Ref. [15] that PL intensity of irradiated samples is practically the same at the Xe ion fluencies of  $10^{12}$  and  $3 \times 10^{12}$   $\text{cm}^{-2}$ . Further the PL signal intensity increases continuously with fluence up to  $10^{14}$   $\text{cm}^{-2}$ . It is correlated with our results. Though, because we covered the more extended fluence range, the PL signal saturation at  $5 \times 10^{14}$   $\text{cm}^{-2}$  was observed.

It is interesting to compare the fluence dependencies of PL intensity and etch velocity (Figs 2 and 4). In general, both PL intensity and etch velocity increases with ion fluence. However, despite similarity between results some differences are observed, too. Unlike of the PL intensity, the etch velocity is continuously increases with fluence up to  $5 \times 10^{14}$   $\text{cm}^{-2}$ , more slowly at low fluencies and more quickly in the fluence range of ( $1 \times 10^{14}$ – $5 \times 10^{14}$ )  $\text{cm}^{-2}$ . One would expect a saturation of the etching velocity at these high fluences because multiple overlapping of individual damaged zones. But, it is not happened. It is known that the differences in the etch velocity are induced by the deformation fields around damaged zones (see, for example [35]). Obviously, the process of mechanical stress accumulation lasts yet at the fluence higher than  $10^{14}$   $\text{cm}^{-2}$  resulting in the modified etching rate. The similar feature of fluence dependencies is practically full recovery of both PL intensity and etch velocity to the initial values after high-temperature treatment.

It is known that the long-time heat treatment results in the decomposition of defect's complexes or their conversion to non-radiative recombination centers. That is why we did not register a noticeable visible luminescence from the high-fluence irradiated samples after annealing at 1100 °C. Tyschenko [36] reported the similar results on the visible PL of a-SiO<sub>2</sub> films implanted with 200 and 100 keV-Si ions to fluencies of  $3 \times 10^{16}$  and  $1.8 \times 10^{16}$   $\text{cm}^{-2}$ , respectively. SiO<sub>2</sub> films were annealed at (400–1000 °C) for 30 min after the irradiation. On the base of analysis of the PL dependence on annealing temperature, the author of Ref [36] concluded that the implantation-induced PL of SiO<sub>2</sub> was due to radiative oxygen deficient centers and noticed that the heat-treatment of implanted SiO<sub>2</sub> at 1000 °C resulted in visible PL quenching. These results are in a good agreement with our results.

As it was mentioned above, a rearrangement of six-component rings consisting of SiO<sub>4</sub>-tetrahedra into three- and four-component rings takes place along the latent SHI track in a-SiO<sub>2</sub> resulting in the emergence of microscopic defects with stressed Si–O bonds in the form of paramagnetic E' centers and oxygen deficient defects [16,25–27]. Heat-treatment is accompanied by a damage recovery in the irradiated SiO<sub>2</sub> film, and the reconstruction of six-component rings from small rings induced by SHIs irradiation takes place again. It should result in the  $V_{irr}$  recovery to the initial value as we observe in our experiment. It should be mentioned that Awazu et al. [16] have studied the structure of latent tracks created in amorphous SiO<sub>2</sub> using the IR absorption spectroscopy and chemical etching in HF. They reported the complete etch rate recovery to the initial value after annealing the irradiated samples at 600 °C for 150 min in inert ambient.

#### 4. Conclusions

We have studied the etch velocity in 4% - HF and the visible PL of 1  $\mu\text{m}$ -thick SiO<sub>2</sub> films thermally grown on Si and afterwards irradiated with 133 MeV Xe ions to the fluencies ranging between

$10^{10}$   $\text{cm}^{-2}$  and  $5 \times 10^{14}$   $\text{cm}^{-2}$ .

It was found a drastic increase of the SiO<sub>2</sub> etch velocity in HF solution at the fluence of  $\geq 10^{14}$   $\text{cm}^{-2}$ . The etch velocity ratio of irradiated and virgin SiO<sub>2</sub> ( $V_{irr}/V_{virgin}$ ) amounts to 2.2 and 4.1 for the fluence of  $1 \times 10^{14}$   $\text{cm}^{-2}$  and  $5 \times 10^{14}$   $\text{cm}^{-2}$ , respectively. Annealing of irradiated samples (1100 °C, 120 min) resulted in a restoration of the etch velocity practically to its virgin value. The  $V_{irr}/V_{virgin}$  ratio can be useful in order to estimate the radiation damage in a-SiO<sub>2</sub> matrix irradiated with high fluencies of SHIs.

An intensive visible luminescence was registered from the samples irradiated to the fluence of  $10^{14}$   $\text{cm}^{-2}$  and higher. The PL signal intensity was saturated at the fluence of  $5 \times 10^{14}$   $\text{cm}^{-2}$ . Kachurin et al. [14,15] consider Si inclusions to be responsible for visible luminescence of SHI-irradiated a-SiO<sub>2</sub>. Presumably, such inclusions can be formed via SiO<sub>2</sub> disproportionation in the ion tracks. In our experiment, TEM and RS data did not prove the formation of Si nanoprecipitates in the as-irradiated a-SiO<sub>2</sub> as well as in the irradiated a-SiO<sub>2</sub> after heat treatment. We believe that an appearance of visible PL is due to the generation of oxygen deficient defects and its complexes in the SiO<sub>2</sub> matrix irradiated with SHIs in conditions of multiple ion track overlapping. The PL signal quenching after annealing of irradiated samples (1100 °C, 120 min) supports additionally a suggestion of SHIs-induced defect's responsibility for a-SiO<sub>2</sub> visible luminescence.

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